Ribbon Operators in 2D Topologically Ordered Spin Systems

Jacob C Bridgeman, Steven T Flammia, David Poulin

ARC Centre of Excellence in Engineered Quantum Systems The University of Sydney

12th November 2015





Overview

• Understanding quantum matter

- 2 dimensional topologically ordered models
- Numerically very challenging: contracting PEPS is PostBQP hard (PostBQP contains QMA, NP, etc.)



Overview

• Understanding quantum matter

- 2 dimensional topologically ordered models
- Numerically very challenging: contracting PEPS is PostBQP hard (PostBQP contains QMA, NP, etc.)

• Exotic phases, disordered in the conventional sense

- No local order parameter, no symmetry
- New exchange statistics not observed in nature



Overview

• Understanding quantum matter

- 2 dimensional topologically ordered models
- Numerically very challenging: contracting PEPS is PostBQP hard (PostBQP contains QMA, NP, etc.)

• Exotic phases, disordered in the conventional sense

- No local order parameter, no symmetry
- New exchange statistics not observed in nature
- Topological codes
 - Quantum information protected against arbitrary errors





 $H = -J \sum_{\langle i,j \rangle} Z_i Z_j$



 $H = -J \sum_{\langle i,j \rangle} Z_i Z_j$



 $H = -J \sum_{\langle i,j \rangle} Z_i Z_j$



$$H = -J\sum_{\langle i,j\rangle} Z_i Z_j - h\sum_j X_j$$



$$H = -J\sum_{\langle i,j\rangle} Z_i Z_j - h\sum_j X_j$$

• Quantum phases: $T \rightarrow 0$ limit

• Only ground states have nonzero Boltzmann weight



$$H = -J\sum_{\langle i,j\rangle} Z_i Z_j - h\sum_j X_j$$

- Quantum phases: $T \rightarrow 0$ limit
 - Only ground states have nonzero Boltzmann weight
- Magnetisation $\frac{1}{N}\sum \langle Z_j \rangle$
 - Can detect ordering with local measurement



 $H = -J \sum_{\langle i,j \rangle} Z_i Z_j$



 $H = -J \sum_{\langle i,j \rangle} Z_i Z_j$



 $H = -J \sum_{\langle i,j \rangle} Z_i Z_j$



 $H = -J \sum_{\langle i,j \rangle} Z_i Z_j$



 $H = -J \sum_{\langle i,j \rangle} Z_i Z_j$



 $H = -J \sum_{\langle i,j \rangle} Z_i Z_j$



 $H = -J \sum_{\langle i,j \rangle} Z_i Z_j$



 $H = -J \sum_{\langle i,j \rangle} Z_i Z_j$



 $H = -J \sum_{\langle i,j \rangle} Z_i Z_j$



 $H = -J \sum_{\langle i,j \rangle} Z_i Z_j$



 $H = -J \sum_{\langle i,j \rangle} Z_i Z_j$



 $H = -J \sum_{\langle i,j \rangle} Z_i Z_j$



 $H = -J \sum_{\langle i,j \rangle} Z_i Z_j$



$$H = -\sum_{\text{blue}} X \frac{X}{X} X - \sum_{\text{green}} Z \frac{Z}{Z} Z$$



$$H = -\sum_{\text{blue}} X \frac{X}{X} X - \sum_{\text{green}} Z \frac{Z}{Z} Z$$

• Every term commutes with every other



$$H = -\sum_{\text{blue}} X \frac{X}{X} X - \sum_{\text{green}} Z \frac{Z}{Z} Z$$

- Every term commutes with every other
- Exactly solvable a good playground



$$H = -\sum_{\text{blue}} X \frac{X}{X} X - \sum_{\text{green}} Z \frac{Z}{Z} Z$$

- Every term commutes with every other
- Exactly solvable a good playground
- Supports many of the characteristic properties of topological ordered models



$$H = -\sum_{\text{blue}} X \frac{X}{X} X - \sum_{\text{green}} Z \frac{Z}{Z} Z$$



$$H = -\sum_{\text{blue}} X \frac{X}{X} X - \sum_{\text{green}} Z \frac{Z}{Z} Z$$



$$H = -\sum_{\text{blue}} X \frac{X}{X} X - \sum_{\text{green}} Z \frac{Z}{Z} Z$$



$$H = -\sum_{\text{blue}} X \frac{X}{X} X - \sum_{\text{green}} Z \frac{Z}{Z} Z$$



$$H = -\sum_{\text{blue}} X \frac{X}{X} X - \sum_{\text{green}} Z \frac{Z}{Z} Z$$



$$H = -\sum_{\text{blue}} X \frac{X}{X} X - \sum_{\text{green}} Z \frac{Z}{Z} Z$$



$$H = -\sum_{\text{blue}} X \frac{X}{X} X - \sum_{\text{green}} Z \frac{Z}{Z} Z$$



$$H = -\sum_{\text{blue}} X \frac{X}{X} X - \sum_{\text{green}} Z \frac{Z}{Z} Z$$



$$H = -\sum_{\text{blue}} X \frac{X}{X} X - \sum_{\text{green}} Z \frac{Z}{Z} Z$$



$$H = -\sum_{\text{blue}} X \frac{X}{X} X - \sum_{\text{green}} Z \frac{Z}{Z} Z$$



$$H = -\sum_{\text{blue}} X \frac{X}{X} X - \sum_{\text{green}} Z \frac{Z}{Z} Z$$



$$H = -\sum_{\text{blue}} X \frac{X}{X} X - \sum_{\text{green}} Z \frac{Z}{Z} Z$$

















• Stringlike = ribbon operators

- Stringlike = ribbon operators
- Commute with the Hamiltonian in the bulk
 - Quasiparticle excitations localised around the ends

- Stringlike = ribbon operators
- Commute with the Hamiltonian in the bulk
 - Quasiparticle excitations localised around the ends
- Exotic exchange statistics



- Stringlike = ribbon operators
- Commute with the Hamiltonian in the bulk
 - Quasiparticle excitations localised around the ends
- Exotic exchange statistics



• Should not depend too much on the specific support chosen

Deformable

• Cost function \implies numerical optimisation

• Cost function \implies numerical optimisation

•
$$C(R;\eta) = \frac{1}{\operatorname{supp} R} \left(\|[R,H]\|^2 + \|RL - \eta LR\|^2 \right)$$

- Cost function \implies numerical optimisation

•
$$C(R;\eta) = \frac{1}{\operatorname{supp} R} \left(\|[R,H]\|^2 + \|RL - \eta LR\|^2 \right)$$

- Search space: minimally entangling operators
 - Matrix product operators of constant bond dimension

- Cost function \implies numerical optimisation

•
$$C(R;\eta) = \frac{1}{\operatorname{supp} R} \left(\|[R,H]\|^2 + \|RL - \eta LR\|^2 \right)$$

- Search space: minimally entangling operators
 - Matrix product operators of constant bond dimension
- Algorithm
 - (Modified) DMRG

Observations: Checking things are working



Observations: Checking things are working

$$H = JH_{\mathbb{Z}_N}$$
 Toric Code $-\frac{h}{2}\sum(Z_j + Z_j^{\dagger}) - \frac{\lambda}{4}\sum_{\langle j,k \rangle}(X_j + X_j^{\dagger})(X_k + X_k^{\dagger})$



Toric Code with Z Field





The Kitaev Honeycomb Model

Frustrated model

• Noncommuting Hamiltonian, symmetries are tricky

The Kitaev Honeycomb Model

Frustrated model

- Noncommuting Hamiltonian, symmetries are tricky
- Known to be in the same phase as \mathbb{Z}_2 Toric Code
 - Form of ribbon operators not known

The Kitaev Honeycomb Model

Frustrated model

- Noncommuting Hamiltonian, symmetries are tricky
- Known to be in the same phase as \mathbb{Z}_2 Toric Code

Form of ribbon operators not known





Review

- New technique for identifying topological order
 - Operator first approach reduces dimensionality







- New technique for identifying topological order
 - Operator first approach reduces dimensionality
- Can be applied to models which cannot be solved analytically
 - Identifying topological order in more realistic models







- New technique for identifying topological order
 - Operator first approach reduces dimensionality
- Can be applied to models which cannot be solved analytically
 - Identifying topological order in more realistic models
- Seems to work for Abelian topological order
 - How to extend to the non-Abelian case? We have some ideas







- New technique for identifying topological order
 - Operator first approach reduces dimensionality
- Can be applied to models which cannot be solved analytically
 - Identifying topological order in more realistic models
- Seems to work for Abelian topological order
 - How to extend to the non-Abelian case? We have some ideas
- Can we prove anything about the method?
 - If we restrict to unitaries, we can prove ground state degeneracy using numerical output





$$H = -J \sum_{i,j} (X_{i,j} X_{i,j+1} + Z_{i,j} Z_{i+1,j})$$



Topologically Encoded Qubit: Quantum Compass Model/Bacon Shor Code

$$H = -J \sum_{i,j} (X_{i,j} X_{i,j+1} + Z_{i,j} Z_{i+1,j})$$





Topologically Encoded Qubit: Quantum Compass Model/Bacon Shor Code

